# Fast Triangular Factorization of the Sum of Quasi-Toeplitz and Quasi-Hankel Matrices* 

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#### Abstract

The literature contains several recent fast algorithms for the triangular factorization of strongly regular Toeplitz-plus-Hankel matrices. In this paper we study the rather more general sum of quasi-Toeplitz and quasi-Hankel matrices, both Hermitian and non-Hermitian. Quasi-Toeplitz and quasi-Hankel matrices are those that are congruent to Toeplitz and Hankel matrices in a special sense. The derivation is based on the concept of displacement structure and its intimate relation to the Schur reduction procedure for triangular factorization. Various special cases covering displacement ranks from two to eight are considered. Several other problems (e.g., factorization of the inverse matrix, solution of exact or overdetermined linear systems) can be reduced to the direct factorization problem.


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## 1. INTRODUCTION

Many signal processing problems require the solution of a linear system of equations of large dimensions. In many applications, (e.g., linear prediction, discrete inverse scattering, cascade synthesis, digital filtering, and the partial realization problem), Toeplitz, quasi-Toeplitz, Hankel, or quasi-Hankel matrices need to be factored or inverted. In many cases the coefficient matrix can be represented as the sum of quasi-Toeplitz and quasi-Hankel matrices. For example, in 1955 Gelfand and Levitan [1] reduced the solution of the inverse scattering problem to the solution of a set of linear integral equations. The discretized version of these equations has a Toeplitz-plus-Hankel coefficient matrix [2]. Such coefficient matrices also arise in the design of least-squares linear-phase prediction and smoothing filters [3] and in the least-squares minimization of forward and backward prediction error energies in $A R$ spectral estimation $[4,5]$. Also, in digital signal processing, one is often faced with the problem of computing the output noise variance due to fixed-point quantization or input signal quantization. This problem was shown [6] to reduce to the solution of a linear system of equations whose coefficient matrix is the sum of a Toeplitz matrix and a quasi-Hankel matrix.

Several fast inversion algorithms [ $O\left(n^{2}\right)$ ] for Toeplitz-plus-Hankel matrices have been published in recent years. Friedlander et al. used the generalized Levinson-Szegö algorithm [7] and the multichannel Levinson algorithm [8] to derive efficient inversion recursions for sums of products of Toeplitz and Hankel matrices [9]. Merchant and Parks [10] reduced the solution of a system of equations with a Toeplitz-plus-Hankel coefficient matrix to the solution of a system of equations with a $2 \times 2$ block Toeplitz coefficient matrix, which is then solved by applying the block Levinson algorithm [11]. This procedure however, does not work for any strongly regular Toeplitz-plus-Hankel matrix, since it requires also the strong regularity of the corresponding Toeplitz-minus-Hankel matrix. Heinig, Jankowski, and Rost [12] also developed fast algorithms for the inversion of strongly regular Toeplitz-plus-Hankel matrices. Their recursive procedure is based on the concept of UV reduction and the solution of the so-called fundamental equations [13]. In [14, 15] Heinig and Rost represented Toeplitz-plus-Hankel inverses as sums of products of triangular Toeplitz and Hankel matrices, and in [16] Lev-Ari used the Schur reduction procedure to derive a fast algorithm for the triangular factorization of Hermitian Toeplitz-plus-Hankel matrices. Gohberg et al. [17] presented a fast algorithm for the solution of a system of linear equations with a strongly regular symmetric coefficient matrix which is the sum of a real Toeplitz matrix and a real Hankel matrix. Their derivation is based on the general approach developed in [17]. More recently, Yagle [18]
extended the split Levinson and Schur algorithms [19, 20] to strongly regular Toeplitz plus Hankel matrices, and Zarowski [21] used the algorithm of Heinig, Jankowski, and Rost [12] to induce a Schur-type recursion for Toeplitz-plus-Hankel matrices.

We remark that many earlier solutions to the factorization of Toeplitz-plus-Hankel matrices have been obtained by embedding the original matrix into a larger Toeplitz matrix. This approach, however, does not exhibit any nesting properties and so can only lead to fixed-order solutions. That is, if the matrix in question grows by one row and one column, then the factorization of the larger matrix has to be obtained by repeating the embedding procedure. In contrast, the algorithm presented in this paper allows for nested solutions: it can easily "update" the previous solution to reflect the change in the original matrix.

In this paper we give fast (and parallelizable) triangular factorization algorithms for a general class of strongly regular structured matrices which can be written as sums of quasi-Toeplitz and quasi-Hankel matrices. This class includes, among others, Toeplitz matrices, Hankel matrices, Toeplitz-plus-Hankel matrices, and the inverse of a Toeplitz matrix plus a Hankel matrix. Both the Hermitian and the non-Hermitian cases are considered. Our derivation is based on the observation that a quasi-Toeplitz matrix plus a quasi-Hankel matrix is structured, though in a more general sense than the "original" notion of displacement structure [22, 23], contrary to the statement made in [21]. We derive a general recursive algorithm (Algorithm 1 in Section 5.1) and then consider some important special cases such as Toeplitz plus Hankel matrices. We also exhibit the recursions in matrix form.

This paper is organized as follows. In Section 2 we review the definition of generating functions, quasi-Toeplitz, quasi-Hankel, and structured matrices, and introduce the family of matrices close to Toeplitz plus Hankel. In Section 3 we show how the Schur reduction procedure yields the triangular factorization of non-Hermitian matrices, and in Section 4 we derive fast triangular factorization algorithms for structured non-Hermitian matrices. These algorithms reduce in the Hermitian case to those already presented in [23]. In Section 5 the derived algorithms are applied to matrices close to Toeplitz plus Hankel, and in Section 6 we specialize these algorithms for quasi-Hankel and quasi-Toeplitz matrices.

## 2. PRELIMINARIES

In this section we review the definition of generating functions, quasiToeplitz matrices, quasi-Hankel matrices, and structured matrices, and intro-
duce the class of matrices close to Toeplitz plus Hankel. Our goal will be triangular factorization, which is a nested operation. Therefore we can assume, without loss of generality, that all matrices are extended to be semiinfinite.

Matrices $T$ that are constant along the diagonals are called Toeplitz matrices. That is, the $(i, j)$ th element is a function of $(i-j)$ :

$$
T=\left[\begin{array}{cccc}
c_{0} & c_{-1} & c_{-2} & \cdots \\
c_{1} & c_{0} & c_{-1} & \\
c_{2} & c_{1} & c_{0} & \\
\vdots & & & \ddots
\end{array}\right]
$$

Similarly, matrices $H$ that are constant along the antidiagonals are called Hankel matrices. That is, the $(i, j)$ th element is a function of $(i+j)$ :

$$
H=\left[\begin{array}{cccc}
h_{0} & h_{1} & h_{2} & \cdots \\
h_{1} & h_{2} & h_{3} & \\
h_{2} & h_{3} & h_{4} & \\
\vdots & & & \ddots
\end{array}\right]
$$

Note that a Hankel matrix $H$ is always symmetric. $H$ is Hermitian if, and only if, $H$ is real, whereas $T$ is Hermitian if, and only if, $c_{-i}=c_{i}^{*}$.

Definition 1 (Cenerating function [23]). The generating function of a semiinfinite matrix $R=\left[r_{i j}\right]_{i, j=0}^{\infty}$ is the bivariate function $R(z, w)$ defined by

$$
R(z, w)=\left[\begin{array}{llll}
1 & z & z^{2} & \ldots
\end{array}\right] R\left[\begin{array}{llll}
1 & w & w^{2} & \ldots
\end{array}\right]^{*}
$$

It can be easily checked that the generating functions of Toeplitz and Hankel matrices are

$$
\begin{equation*}
T(z, w)=\frac{c(z)+r^{*}(w)}{1-z w^{*}} \quad \text { and } \quad H(z, w)=\frac{j\left[z h(z)-w^{*} h\left(w^{*}\right)\right]}{j\left(z-w^{*}\right)} \tag{1}
\end{equation*}
$$

where $j=\sqrt{-1}$ and

$$
\begin{aligned}
& h(z)=h_{0}+z h_{1}+z^{2} h_{2}+\cdots \\
& c(z)=\frac{c_{0}}{2}+z c_{1}+z^{2} c_{2}+\cdots \\
& r^{*}(w)=\frac{c_{0}}{2}+w^{*} c_{-1}+w^{* 2} c_{-2}+\cdots
\end{aligned}
$$

Notice that when $H$ is Hermitian (or real), $h\left(z^{*}\right)=h^{*}(z)$.

Defintion 2 (Quasi-Toeplitz matrix [23]). A matrix $Q T$ is said to be quasi-Toeplitz if its generating function $\mathrm{QT}(z, w)$ can be written in the form

$$
\begin{equation*}
\operatorname{QT}(z, w)=a(z) T(z, w) b^{*}(w) \tag{2}
\end{equation*}
$$

for some Toeplitz matrix $T$ and functions $a(z)$ and $b(z)$. The matrix QT is Hermitian if $a(z)=b(z)$ and $T$ is Hermitian.

Definition 3 (Quasi-Hankel matrix [23]). A matrix QH is said to be quasi-Ilankel if its generating function $\operatorname{QII}(z, w)$ can be written in the form

$$
\begin{equation*}
Q \mathrm{H}(z, w)=\alpha(z) H(z, w) \beta^{*}(w) \tag{3}
\end{equation*}
$$

for some Hankel matrix $H$ and functions $\alpha(z)$ and $\beta(z)$. The matrix QH is Hermitian if $\alpha(z)=\beta(z)$ and $H$ is real.

In this paper, we shall study matrices $R$ whose generating functions can be written in the form

$$
R(z, w)=\mathrm{QT}(z, w)+\mathrm{QH}(z, w) .
$$

That is, $R$ is the sum of a quasi-Toeplitz matrix and a quasi-Hankel matrix. This clearly includes the sum of Toeplitz and Hankel matrices ( $R=T+H$ ). It also includes matrices $R$ of the form $R=T^{-1}+H$, since the inverse of a Toeplitz matrix is quasi-Toeplitz (see [23] and the references therein).

Using the expression (1) for $T(z, w)$ and $H(z, w)$, and the expressions (2)-(3), we obtain that $R(z, w)$ is of the form

$$
\begin{equation*}
R(z, w)=\frac{G(z) J B^{*}(w)}{j\left(1-z w^{*}\right)\left(z-w^{*}\right)}, \tag{4}
\end{equation*}
$$

where $G(z)$ and $B(z)$ are $1 \times 8$ row matrix functions

$$
\begin{aligned}
G(z) & =\left[\begin{array}{lllll}
a(z) & z a(z) & \alpha(z) & z \alpha(z) & z^{2} \alpha(z) h(z) \\
& z \alpha(z) h(z) & a(z) c(z) & z a(z) c(z)
\end{array}\right], \\
B(z) & =\left[\begin{array}{lllll}
b(z) & z b(z) & \beta(z) & z \beta(z) & z^{2} \beta(z) \tilde{h}(z) \\
& z \beta(z) \tilde{h}(z) & b(z) r(z) & z b(z) r(z)
\end{array}\right], \\
J & =\left[\begin{array}{llll} 
\\
j & j & -j &
\end{array}\right]
\end{aligned}
$$

and $\tilde{h}(z)=\sum_{i=0}^{\infty} h_{i}^{*} z^{i}$. Notice that $\tilde{h}(z)=h^{*}\left(z^{*}\right)$ is the same power series as $h(z)$ but with conjugate coefficients. If $a(z)=b(z), \alpha(z)=\beta(z)$, and $H$ is real, then $\breve{h}(z)=h(z)$ and $G(z)=B(z)$, and thus $R$ is Hermitian. In some special cases, the matrices $G(z)$ and $B(z)$ can have a lower column dimension. For example, if $a(z)=\alpha(z)$ and $b(z)=\beta(z)$, then

$$
\left.\left.\left.\left.\begin{array}{l}
G(z)=\left[\begin{array}{llll}
a(z) & z a(z) & a(z)\left[c(z)+z^{2} h(z)\right.
\end{array}\right] \quad z a(z)[c(z)+h(z)
\end{array}\right]\right], \text {, } \begin{array}{lll}
b(z) & z b(z) & b(z)\left[r(z)+z^{2} \tilde{h}(z)\right.
\end{array}\right] \quad z b(z)[r(z)+\tilde{h}(z)]\right], ~ \$\left[\begin{array}{lll}
b(z) \\
B(z)
\end{array}\right.
$$

and

$$
J=\left[\begin{array}{llll} 
& & & -j \\
& & j & \\
& -j & &
\end{array}\right]
$$

This special case includes Toeplitz-plus-Hankel matrices $[a(z)=b(z)=$ $\alpha(z)=\beta(z)=1$ ].

Equation (4) and its special cases belong to the more general family whose generating functions can be written as

$$
\begin{equation*}
R(z, w)=\frac{G(z) J B^{*}(w)}{d(z, w)} \tag{5}
\end{equation*}
$$

where $J$ is any constant nonsingular matrix, $d(z, w)$ is the generating function of a constant (possibly singular) Hermitian matrix $\boldsymbol{d}=\left[d_{i j}\right]_{i, j=0}^{\infty}$, viz.,

$$
d(z, w)=\left[\begin{array}{llll}
1 & z & z^{2} & \cdots
\end{array}\right] \boldsymbol{d}\left[\begin{array}{llll}
1 & w & w^{2} & \cdots
\end{array}\right]^{*}
$$

and $G(z)$ and $B(z)$ are $1 \times \rho$ row vector functions, where $\rho$ is called the displacement rank of $R$. Equation (5) has a matrix-domain equivalent. If we define

$$
\begin{equation*}
\nabla_{d} R=\sum_{i, j=0}^{\infty} d_{i j} Z^{i} R Z^{* j} \tag{6}
\end{equation*}
$$

where Z is the lower triangular shift matrix with ones on the first subdiagonal, then

$$
\nabla_{d} R=G J B^{*} \quad \text { with } \quad \rho=\operatorname{rank}\left(\nabla_{d} R\right)
$$

where $G$ and $B$ are matrices with $\rho$ columns. Each column of $G$ [respectively $B]$ is formed by stacking the coefficients of the corresponding function in $G(z)$ [respectively $B(z)$ ], i.e.,

$$
G(z)=\left[\begin{array}{llll}
1 & z & z^{2} & \ldots
\end{array}\right] G \quad \text { and } \quad B(z)=\left[\begin{array}{llll}
1 & z & z^{2} & \cdots \tag{7}
\end{array}\right] B
$$

The structure (5) allows a so-called fast Schur reduction procedure for triangular factorization. This is described in the next two sections.

## 3. THE SCHUR REDUCTION PROCEDURE

The Schur reduction is an $O\left(n^{3}\right)$ procedure that factors an $n \times n$ matrix $R$ into the form

$$
R=L D U
$$

where $L$ is a lower triangular matrix with unit diagonal elements, $U$ is an upper triangular matrix with unit diagonal elements, and $D$ is a diagonal matrix. A necessary and sufficient condition for $D$ to be nonsingular is that $R$ be strongly regular (i.e., all leading principal submatrices of $R$ are nonsingular). This condition will be assumed throughout this paper. For a strongly regular matrix $R=\left[r_{i j}\right]_{i, j=0}^{n-1}$ we have

$$
R-\left[\begin{array}{c}
r_{00} \\
r_{10} \\
r_{20} \\
\vdots
\end{array}\right] r_{00}^{-1}\left[\begin{array}{llll}
r_{00} & r_{01} & r_{02} & \cdots
\end{array}\right]=\left[\begin{array}{ccccc}
0 & 0 & \cdots & 0 & 0 \\
0 & & & & \\
\vdots & & \tilde{R}_{1} & & \\
0 & & & &
\end{array}\right]
$$

Let $d_{0}=r_{00}$, and define
$l_{0}=\left[\begin{array}{c}r_{00} \\ r_{10} \\ r_{20} \\ \vdots\end{array}\right] d_{0}^{-1}, \quad u_{0}=\left[\begin{array}{c}r_{00}^{*} \\ r_{01}^{*} \\ r_{02}^{*} \\ \vdots\end{array}\right] d_{0}^{-*}, \quad$ and $\quad R_{1}=\left[\begin{array}{ccccc}0 & 0 & \cdots & 0 & 0 \\ 0 & & & & \\ \vdots & & \tilde{R}_{1} & & \\ 0 & & & & \end{array}\right] ;$
then $R=d_{0} l_{0} u_{0}^{*}+R_{1}$. The matrix $\tilde{R}_{1}$ is called the Schur complement of $R$ with respect to $r_{00}$. This suggests the following recursive procedure for the triangular factorization of $R$ :

$$
\begin{gather*}
R_{0}=R, \quad R_{i+1}=R_{i}-d_{i} l_{i} u_{i}^{*}, \quad 0 \leqslant i \leqslant n-1 \\
d_{i}-e_{i}^{T} R_{i} e_{i}, \quad l_{i}=R_{i} e_{i} d_{i}^{-1}, \quad \text { and } \quad u_{i}^{*}=e_{i}^{T} R_{i} d_{i}^{-1} \tag{8}
\end{gather*}
$$

where $e_{i}=\left[\begin{array}{lllllll}0 & \cdots & 0 & 1 & 0 & \cdots & 0\end{array}\right]^{T}$ has $i$ leading zeros. Notice that the first $i$
rows and columns of $R_{i}$ are zero, the first $i$ elements of $l_{i}$ and $u_{i}$ are zero, and the $(i+1)$ th elements of $l_{i}$ and $u_{i}$ are both equal to 1 . For an $n \times n$ matrix $R$, the matrices $L, D$, and $U$ are given by

$$
\begin{gathered}
L=\left[\begin{array}{llll}
l_{0} & l_{1} & \cdots & l_{n-1}
\end{array}\right], \quad U=\left[\begin{array}{llll}
u_{0} & u_{1} & \cdots & u_{n-1}
\end{array}\right]^{*}, \\
D=\operatorname{diag}\left\{d_{0}, d_{1}, \ldots, d_{n-1}\right\} .
\end{gathered}
$$

The diagonal elements $d_{i}$ are guaranteed to be nonzero by the strong regularity of $R$. For a matrix $R$ that is not strongly regular the recursive procedure breaks down if either or both of the following situations arise:

$$
\begin{array}{ll}
d_{i}=0 & \text { and } \quad R_{i} e_{i} \neq \mathbf{0} \\
d_{i}=0 & \text { and } \quad e_{i}^{T} \boldsymbol{R}_{i} \neq \mathbf{0}
\end{array}
$$

However, if $d_{i}=0, R_{i} e_{i}=\mathbf{0}$, and $e_{i}^{T} R_{i}=\mathbf{0}$, then we can choose $l_{i}$ and $u_{i}$ arbitrarily. A convenient choice would be $l_{i}=u_{i}=e_{i}$, which is consistent with $L$ and $U$ having unit diagonal elements.

## 4. GENERATING FUNCTION APPROACH

The complexity of the Schur reduction procedure can be reduced to $O\left(\rho n^{2}\right)$ in the case of structured matrices as in (5). This is because the recursive procedure (8) reduces to a recursive update of the generator matrices $G$ and $B$, which have $\rho n$ elements each, as compared to $n^{2}$ in $R$. We extend here the approach used in [23,24] to the non-Hermitian case. Using the notation of generating functions (and assuming, without loss of generality, that $R$ is extended to be semiinfinite as explained at the beginning of Section 2), it is easy to check that (8) can be written in the form

$$
\begin{gather*}
d_{i}=\tilde{R}_{i}(0,0)  \tag{9}\\
l_{i}(z)=z^{i} \tilde{R}_{i}(z, 0) d_{i}^{-1} \quad \text { and } \quad u_{i}(w)=w^{* i} \tilde{R}_{i}(0, w) d_{i}^{-1},  \tag{10}\\
z w^{*} \tilde{R}_{i+1}(z, w)=\tilde{R}_{i}(z, w)-\tilde{R}_{i}(z, 0) \tilde{R}_{i}^{-1}(0,0) \tilde{R}_{i}(0, w), \tag{11}
\end{gather*}
$$

where

$$
\begin{aligned}
l_{i}(z) & =\left[\begin{array}{llll}
1 & z & z^{2} & \ldots
\end{array}\right] l_{i} \\
u_{i}(w) & =u_{i}^{*}\left[\begin{array}{llll}
1 & w & w^{2} & \ldots
\end{array}\right]^{*}, \\
\tilde{R}_{i}(z, w) & =\left(z w^{*}\right)^{-i} R_{i}(z, w) .
\end{aligned}
$$

Notice that the triangular factors of the given finite matrix $R$ are obtained by considering the first $n$ coefficients of $l_{i}(z)$ and $u_{i}(z)$.

We shall now proceed by induction. Suppose

$$
\tilde{R}_{i}(z, w)=\frac{G_{i}(z) J B_{i}^{*}(w)}{d(z, w)}
$$

(this is certainly true for $i=0$ ); then (11) yields

$$
z w^{*} \tilde{R}_{i+1}(z, w)=\frac{G_{i}(z)\left(J-\frac{d(z, w)}{d(z, 0) d(0, w)} J M_{i} J\right) B_{i}^{*}(w)}{d(z, w)}
$$

where

$$
M_{i}=B_{i}^{*}(0) \tilde{R}_{i}^{-1}(0,0) G_{i}(0)
$$

Observe that $M_{i}$ satisfies

$$
\begin{equation*}
M_{i} J M_{i}=d(0,0) M_{i} \tag{12}
\end{equation*}
$$

If we can find matrices $\Theta_{i}(z)$ and $\Gamma_{i}(z)$ such that

$$
\begin{equation*}
\Theta_{i}(z) J \Gamma_{i}^{*}(w)=J-\frac{d(z, w)}{d(z, 0) d(0, w)} J M_{i} J \tag{13}
\end{equation*}
$$

then we can write

$$
z w^{*} \tilde{R}_{i+1}(z, w)=\frac{G_{i}(z) \Theta_{i}(z) J \Gamma_{i}^{*}(w) B_{i}^{*}(w)}{d(z, w)}
$$

This shows that we can use the following recursions:

$$
\begin{array}{ll}
z G_{i+1}(z)=G_{i}(z) \Theta_{i}(z), & G_{0}(z)=G(z) \\
z B_{i+1}(z)=B_{i}(z) \Gamma_{i}(z), & B_{0}(z)=B(z) \tag{14}
\end{array}
$$

for the update of $G_{i}(z)$ and $B_{i}(z)$. Moreover, $d_{i}, l_{i}(z)$, and $u_{i}(z)$ can be determined from $G_{i}(z)$ and $B_{i}(z)$ without the need to evaluate $\vec{R}_{i}(z, w)$ explicitly:

$$
\begin{gathered}
d_{i}=\lim _{z \rightarrow 0} \frac{G_{i}(z) J B_{i}^{*}(z)}{d(z, z)} \\
l_{i}(z)=z^{i} \frac{G_{i}(z) J B_{i}^{*}(0)}{d(z, 0)} d_{i}^{-1}, \quad \text { and } \quad u_{i}(w)=w^{* i} \frac{G_{i}(0) J B_{i}^{*}(w)}{d(0, w)} d_{i}^{-1}
\end{gathered}
$$

Since the matrices $G_{i}$ and $B_{i}$ have each $\rho(n-i)$ elements, then $n$ steps of the recursions (14) require $O\left(\rho n^{2}\right)$ operations (additions and multiplications), which represents great savings in computation when $\rho<n$.

The existence of solutions $\Theta_{i}(z)$ and $\Gamma_{i}(z)$ to Equation (13) is guaranteed by the following theorem.

THEOREM 1 (Factorization). Let $R$ be a strongly regular structured matrix whose generating function is given by

$$
R(z, w)=\frac{G(z) J B^{*}(w)}{d(z, w)}
$$

where $J$ is a constant nonsingular matrix and $d(z, w)$ is a scalar bivariate Hermitian function of the form

$$
\begin{equation*}
d(z, w)=e(z) e^{*}(w)-f(z) f^{*}(w) \tag{15}
\end{equation*}
$$

Then the triangular factorization of $R$ can be carried out recursively using (14) with

$$
\Theta_{i}(z)=\left\{I-\lambda(z) J M_{i}\right\} U_{i} \quad \text { and } \quad \Gamma_{i}(z)=\left\{I-\lambda(z) J^{*} M_{i}^{*}\right\} V_{i}
$$

where

$$
\begin{equation*}
\lambda(z)=\frac{d(z, \tau)}{d(z, 0) d(0, \tau)} \tag{16}
\end{equation*}
$$

$U_{i}$ and $V_{i}$ are arbitrary constant matrices satisfying $U_{i} J V_{i}^{*}=J$, and $\tau$ is any scalar such that $d(\tau, \tau)=0$.

Proof. This theorem is an extension to the non-Hermitian case of Theorem 3 in [23] and of the discussion in [24]. Using the expressions for $\Theta_{i}(z)$ and $\Gamma_{i}(z)$ with (12), we have

$$
\Theta_{i}(z) J \Gamma_{i}^{*}(w)=J-\left[\lambda(z)+\lambda^{*}(w)-\lambda(z) \lambda^{*}(w) d(0,0)\right] J M_{i} J
$$

Substituting into (13), we obtain

$$
\frac{d(z, w)}{d(z, 0) d(0, w)}=\lambda(z)+\lambda^{*}(w)-d(0,0) \lambda(z) \lambda^{*}(w)
$$

This is the same as an equation obtained in [23, 24], where it is shown that a solution $\lambda(z)$ exists if, and only if, $d(z, w)$ is of the form (15). Moreover, $\lambda(z)$ is given by (16).

Corollary 1 (Generators of the Schur complement). $G_{i}$ and $B_{i}$ are the generator matrices of $\tilde{R}_{i}$, which is the Schur complement of $R$ with respect to the leading $i \times i$ principal submatrix.

## 5. FACTORIZATION OF QT + QH MATRICES

In this section we specialize the recursions of the previous theorem for strongly regular matrices $R$ whose generating functions can be expressed in the form

$$
\begin{equation*}
R(z, w)=\frac{G(z) J B^{*}(w)}{j\left(1-z w^{*}\right)\left(z-w^{*}\right)} \tag{17}
\end{equation*}
$$

where $G$ and $B$ are matrices with $\rho$ columns each, and $J$ is a $\rho \times \rho$ matrix
with nonzero elements along the antidiagonal given by an alternating sequence of $-j$ and $j$ :

$$
J=\left[\begin{array}{llll} 
& & & -j \\
& -j & \\
. & &
\end{array}\right.
$$

It is clear that $J$ is Hermitian when $\rho$ is even, and anti-Hermitian (i.e. $J^{*}=-J$ ) when $\rho$ is odd. We shall denote by $J$ the real matrix defined by $J=j \tilde{J}$.

Next observe that $d(z, w)$ in (17) can be written as

$$
d(z, w)=\left[\begin{array}{ll}
1+z^{2} & z
\end{array}\right]\left[\begin{array}{cc}
0 & -j \\
j & 0
\end{array}\right]\left[\begin{array}{c}
1+w^{* 2} \\
w^{*}
\end{array}\right]
$$

This shows that $d(z, w)$ is in fact the generating function of a rank-2 Hermitian matrix $\mathbf{d}$, which has one positive and one negative eigenvalue and thus is of the form (15). Hence, quasi-Toeplitz-plus-quasi-Hankel matrices are structured with respect to the following matrix-domain displacement (recall (6), and see [25, 26] for generalized definitions of displacement structure):

$$
\nabla_{d} R=Z R\left(I+Z^{2}\right)-\left(I+Z^{2}\right) R Z^{*}
$$

We should remark, however, that the sum of two structured matrices, each one with a different $d(z, w)$, is not in general structured. That is, the resulting $d(z, w)$ need not in general have the form described by (15). Our derivation shows that for the special case of sums of QT and QH matrices, the corresponding QT + QH matrix is structured and hence the factorization procedure described in the previous section can be applied. Observe that this procedure has many degrees of freedom represented by the parameters $U_{i}$, $V_{i}$, and $\tau$. In the sequel we shall make the simplest choices so as to simplify the final algorithms. We have

$$
\lambda(z)=j\left(1-z \tau^{*}\right)\left(\frac{1}{\tau^{*}}-\frac{1}{z}\right),
$$

where $\tau$ is any solution of $\left(1-\tau \tau^{*}\right)\left(\tau-\tau^{*}\right)=0$. The simplest form of $\lambda(z)$ is obtained for $\tau=j$. In this case $\lambda(z)=-j\left(z+z^{-1}\right)$. We are assuming that
$R$ is strongly regular and thus $d_{i}=\tilde{R_{i}}(0,0) \neq 0$ for every $i$. But $d(0,0)$ in (17) is zero, which shows that

$$
\begin{equation*}
G_{i}(0) \tilde{J B_{i}^{*}}(0)=0 \tag{18}
\end{equation*}
$$

Therefore we need a way to compute $\tilde{R}_{i}(0,0)$.
Lemma 1 (Diagonal factors). $\quad \tilde{R}_{i}(0,0)$ can be computed by either of the following two expressions:

$$
\begin{aligned}
& \tilde{R}_{i}(0,0)=G_{i}^{\prime}(0) \tilde{J} B_{i}^{*}(0) \\
& \tilde{R}_{i}(0,0)=-G_{i}(0) \tilde{J} B_{i}^{\prime *}(0)
\end{aligned}
$$

where $G_{i}^{\prime}(0)$ (respectively $B_{i}^{\prime}(0)$ ) is the second row of the corresponding matrix $G_{i}$ (respectively $B_{i}$ ).

Proof. The proof follows easily by noting that the column vectors $l_{i}$ and $u_{i}$ both have $i$ leading zeros followed by 1 . From (8) and (10) we get

$$
1=G_{i}^{\prime}(0) \tilde{J} B_{i}^{*}(0) d_{i}^{-1}=-G_{i}(0) \tilde{J B_{i}^{\prime *}}(0) d_{i}^{-1}
$$

The choice of $U_{i}$ and $V_{i}$ is nonunique. We may choose $U_{i}=V_{i}=I$, which leads to an overall $O\left(4 \rho n^{2}\right)$ computational complexity for each generator. An alternative choice would be to reduce the generators $G_{i}$ and $B_{i}$ to proper form. More specifically, since $J$ is nonsingular and $U_{i} J V_{i}{ }^{*}=J$ (which can be written also as $U_{i} \tilde{J} V_{i}^{*}=\tilde{J}$ ), the matrices $U_{i}$ and $V_{i}$ are nonsingular. Moreover, neither $G_{i}(0)$ nor $B_{i}(0)$ can be 0 , because $d_{i}$ would then be zero (by Lemma 1) and $R$ would not be a strongly regular matrix. Hence it is always possible to find matrices $U_{i}$ and $V_{i}$ such that $G_{i} U_{i}$ and $B_{i} V_{i}$ are reduced to the forms

$$
\begin{gather*}
U_{i} \tilde{J}_{i}^{*}=\tilde{J}  \tag{19}\\
G_{i}(0) U_{i}=\left|\tilde{R}_{i}(0,0)\right|^{1 / 2}\left[\begin{array}{llll}
1 & 0 & \cdots & 0
\end{array}\right]  \tag{20}\\
B_{i}(0) V_{i}=\left|\tilde{R}_{i}(0,0)\right|^{1 / 2}\left[\begin{array}{llll}
1 & 0 & \cdots & 0
\end{array}\right] \tag{21}
\end{gather*}
$$

We can, for instance, use a sequence of elementary rotations to achieve (20)-(21). This also requires $O\left(4 \rho n^{2}\right)$ operations per generator. We shall however give alternative (explicit) expressions for $U_{i}$ and $V_{i}$ for all $\rho \leqslant 8$.

Meanwhile, suppose we find matrices $U_{i}$ and $V_{i}$ that satisfy (19)-(21); then

$$
\begin{align*}
\Theta_{i}(z) & =U_{i}-\left(z+z^{-1}\right) \tilde{J}_{i}^{*}(0) \tilde{R}_{i}^{-1}(0,0) G_{i}(0) U_{i} \\
& =U_{i}\left\{I-\left(z+z^{-1}\right) U_{i}{ }^{1} \tilde{J} B_{i}^{*}(0) \tilde{R}_{i}^{-1}(0,0) G_{i}(0) U_{i}\right\} \\
& =U_{i}\left\{I-\left(z+z^{-1}\right) \tilde{J} V_{i}^{*} B_{i}^{*}(0) \tilde{R}_{i}^{-1}(0,0) G_{i}(0) U_{i}\right\} \\
& =U_{i} S_{\Theta}(z), \tag{22}
\end{align*}
$$

where $S_{\Theta}(z)$ is the $\rho \times \rho$ matrix

$$
S_{\Theta}(z)=I+\left[\begin{array}{cccc}
0 & & 0 & \\
\vdots & & & \\
0 & & & \\
-\epsilon_{i}\left(z+z^{-1}\right) & 0 & \cdots & 0
\end{array}\right]
$$

and

$$
\epsilon_{i}= \begin{cases}+\frac{\left|\tilde{R}_{i}(0,0)\right|}{\tilde{R}_{i}(0,0)} & \text { if } \rho \text { is even }  \tag{23}\\ -\frac{\left|\tilde{R}_{i}(0,0)\right|}{\tilde{R}_{i}(0,0)} & \text { if } \rho \text { is odd }\end{cases}
$$

Similarly we can show that $\Gamma_{i}(z)$ reduces to

$$
\begin{equation*}
\Gamma_{i}(z)=V_{i} S_{\Gamma}(z) \tag{24}
\end{equation*}
$$

where

$$
S_{\Gamma}(z)=I+\left[\begin{array}{cccc}
0 & & 0 & \\
\vdots & & & \\
0 & \text { and } \quad \varepsilon_{i}=\frac{\left|\tilde{R}_{i}(0,0)\right|}{\tilde{R}_{i}^{*}(0,0)} . . . ~ . ~
\end{array}\right.
$$

### 5.1. Matrix Forms

It is easy to verify that the generator recursions (14) [along with (22) and (24)] have the following matrix representations:

$$
\begin{aligned}
& {\left[\begin{array}{c}
\mathbf{0} \\
G_{i+1}
\end{array}\right]=G_{i} U_{i}-\epsilon_{i}\left(Z+Z^{*}\right) G_{i} U_{i}\left[\begin{array}{ll}
0_{\rho-1} \\
1 &
\end{array}\right]} \\
& {\left[\begin{array}{c}
\mathbf{0} \\
B_{i+1}
\end{array}\right]=B_{i} V_{i}-\varepsilon_{i}\left(Z+Z^{*}\right) B_{i} V_{i}\left[\begin{array}{ll}
0_{\rho-1} \\
1 &
\end{array}\right]}
\end{aligned}
$$

where $G_{i}(z)$ and $G_{i}$ [similarly for $B_{i}(z)$ and $B_{i}$ ] are related as in (7). This shows that $G_{i+1}$ is obtained as follows (a similar argument holds for $B_{i+1}$ ):
(a) We multiply $G_{i}$ by $U_{i}$ and keep the last $\rho-1$ columns.
(b) The last column of $G_{i} U_{i}$ is shifted up (introducing a zero at the bottom) and shifted down (introducing a zero at the top). The two shifted versions are added together and multiplied by $-\epsilon_{i}$. The result is then added to the first column of $G_{i} U_{i}$. This gives the first column of $G_{i+1}$.
Moreover, the nonzero parts of the (semiinfinite) triangular factors $l_{i}$ and $u_{i}$ (denoted by $\tilde{l}_{i}$ and $\tilde{u}_{i}$ respectively),

$$
l_{i}=\left[\begin{array}{c}
\mathbf{0}_{i \times 1} \\
\tilde{l_{i}}
\end{array}\right] \quad \text { and } \quad u_{i}=\left[\begin{array}{c}
\mathbf{0}_{i \times 1} \\
\tilde{u}_{i}
\end{array}\right]
$$

are given by

$$
\begin{gathered}
\tilde{l}_{i}(z)=z^{-1} G_{i}(z) \tilde{J} B_{i}^{*}(0) d_{i}^{-1} \\
\tilde{u}_{i}(w)=-w^{-*} G_{i}(0) \tilde{J} B_{i}^{*}(w) d_{i}^{-1}
\end{gathered}
$$

which can be rewritten in matrix form as well. Consider the expression for $\tilde{l}_{i}(z)$ [the same argument holds for $\tilde{u}_{i}(w)$ ]. We write

$$
\begin{aligned}
\tilde{l}_{i} & =Z^{*} G_{i} U_{i} J V_{i}^{*} B_{i}^{*}(0) \tilde{R}_{i}^{-1}(0,0) \\
& =\epsilon_{i}\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} Z^{*} g_{i}
\end{aligned}
$$

where $g_{i}$ denotes the last column of $G_{i} U_{i}$. Recall that $\tilde{l}_{i}$ is a column vector starting with 1 . Hence the second element of $g_{i}$ must be equal to $\left(\epsilon_{i}\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2}\right)^{-1}$. This follows also from Lemma 1. Similarly, we can show that

$$
\tilde{u}_{i}=\varepsilon_{i}\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} Z^{*} b_{i}
$$

where $b_{i}$ is the last column of $B_{i} V_{i}$. Moreover, the second element of $b_{i}$ must be equal to $\left(\varepsilon_{i}\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2}\right)^{-1}$.

In summary, we have the following algorithm.

Algorithm 1 (Factorization of $\mathrm{QT}+\mathrm{QH}$ ). The triangular factorization of a matrix $R$ with generating function $R(z, w)$ as in (17) can be carried out recursively as follows:

$$
\left.\begin{array}{ll}
{\left[\begin{array}{c}
0 \\
G_{i+1}
\end{array}\right]=G_{i} U_{i}-\epsilon_{i}\left(Z+Z^{*}\right) G_{i} U_{i}[ } & 0_{\rho-1} \\
1 &
\end{array}\right], \quad G_{0}=G, ~\left(\begin{array}{c}
\mathbf{0} \\
B_{i+1}
\end{array}\right]=B_{i} V_{i}-\varepsilon_{i}\left(Z+Z^{*}\right) B_{i} V_{i}\left[\begin{array}{cc}
0_{p-1} \\
1 &
\end{array}\right], \quad B_{0}=B .
$$

Let $g_{i}$ and $b_{i}$ denote the last columns of $G_{i} U_{i}$ and $B_{i} V_{i}$ respectively, and let $g_{i 2}$ and $b_{i 2}$ denote the second elements of $g_{i}$ and $b_{i}$ respectively. Then

$$
\tilde{l}_{i}=Z^{*} \frac{g_{i}}{g_{i 2}} \quad \text { and } \quad \tilde{u}_{i}=Z^{*} \frac{b_{i}}{b_{i 2}}
$$

The diagonal factor $d_{i}$ [or $\left.\tilde{R}_{i}(0,0)\right]$ is computed as in Lemma 1, and $\epsilon_{i}$ and $\varepsilon_{i}$ as defined above.

This leaves only the problem of finding $U_{i}$ and $V_{i}$. As we mentioned earlier, there are many possible choices. We shall give explicit expressions for
two typical examples ( $\rho=4,5$ ) and the same approach holds for other cases that we list in Appendix A. Meanwhile, it is worth noting that the recursions in the algorithm get simplified in the following cases:
(a) $\rho$ is even and $G(z)=B(z)$ : in this case $R$ is Hermitian, and it follows from (8) that $l_{i}=u_{i}$ and $d_{i}^{*}=d_{i}$ (i.e., $d_{i}$ is real for all $i$ ).
(b) $\rho$ is odd and $G(z)=B(z)$ : in this case $R$ is anti-Hermitian, and it follows from (8) that $l_{i}=u_{i}$ and $d_{i}^{*}=-d_{i}$ (i.e., $d_{i}$ is imaginary for all $i$ ).

EXAMPLE 1 (Displacement rank 4). This example includes the case of Toeplitz-plus-Hankel matrices. Define the quantities

$$
\begin{aligned}
& {\left[\begin{array}{llll}
\gamma_{i}^{*} & -\delta_{i}^{*} & \eta_{i}^{*} & -\kappa_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} G_{i}(0),} \\
& {\left[\begin{array}{llll}
\mu_{i}^{*} & -\sigma_{i}^{*} & \chi_{i}^{*} & -\psi_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} B_{i}(0),} \\
& U_{i}=\left[\begin{array}{cccc}
\gamma_{i}^{-*} & \delta_{i}^{*} & \delta_{i}^{*}-\frac{\eta_{i}^{*}}{\mu_{i} \gamma_{i}^{*}} & \psi_{i} \\
0 & \gamma_{i}^{*} & \gamma_{i}^{*} & \chi_{i} \\
0 & 0 & \mu_{i}^{-1} & \sigma_{i} \\
0 & 0 & 0 & \mu_{i}
\end{array}\right], \quad \text { and } \\
& V_{i}=\left[\begin{array}{cccc}
\mu_{i}^{-*} & \sigma_{i}^{*} & \sigma_{i}^{*}-\frac{\chi_{i}^{*}}{\gamma_{i} \mu_{i}^{*}} & \kappa_{i} \\
0 & \mu_{i}^{*} & \mu_{i}^{*} & \eta_{i} \\
0 & 0 & \gamma_{i}^{-1} & \delta_{i} \\
0 & 0 & 0 & \gamma_{i}
\end{array}\right] .
\end{aligned}
$$

From (18), it follows that

$$
\begin{equation*}
-\kappa_{i}^{*} \mu_{i}+\eta_{i}^{*} \sigma_{i}-\delta_{i}^{*} \chi_{i}+\gamma_{i}^{*} \psi_{i}=0 \tag{25}
\end{equation*}
$$

It is easy to check with the help of (25) that (19)-(21) are satisfied. Multiplying $G_{i}$ by $U_{i}$ (similarly for $B_{i} V_{i}$ ) requires $O\left(\left(\rho^{2} / 2\right)(n-i)\right.$ ) operations, which leads to an overall $O\left(\left(\rho^{2} / 2\right) n^{2}\right)$ complexity. In this example $\rho=4$; hence we need $O\left(8 n^{2}\right)$ operations, which is of the same order as $O\left(12 n^{2}\right)$ obtained by choosing $U_{i}=V_{i}=I$, or by using a sequence of elementary rotations.

For the case of Toeplitz-plus-Hankel matrices, we start the recursions with

$$
\left.\begin{array}{l}
G(z)=\left[\begin{array}{llll}
1 & z & c(z)+z^{2} h(z) & z[c(z)+h(z)
\end{array}\right]
\end{array}\right],
$$

or equivalently, we start the matrix recursions of Algorithm 1 with the $2 n \times 4$ generators $G$ and $B$ given by

$$
C=\left[\begin{array}{cccc}
1 & 0 & \frac{c_{0}}{2} & 0 \\
0 & 1 & c_{1} & \frac{c_{0}}{2}+h_{0} \\
0 & 0 & c_{2}+h_{0} & c_{1}+h_{1} \\
\vdots & \vdots & \vdots & \vdots \\
0 & 0 & c_{n-1}+h_{n-3} & c_{n-2}+h_{n-2} \\
0 & 0 & h_{n-2} & c_{n-1}+h_{n-1} \\
0 & 0 & h_{n-1} & h_{n} \\
\vdots & \vdots & \vdots & \vdots \\
0 & 0 & h_{2(n-2)} & h_{2(n-1)}
\end{array}\right]
$$

and

$$
B=\left[\begin{array}{cccc}
1 & 0 & \frac{c_{0}}{2} & 0 \\
0 & 1 & c_{-1}^{*} & \frac{c_{0}}{2}+h_{0}^{*} \\
0 & 0 & c_{-2}^{*}+h_{0}^{*} & c_{-1}^{*}+h_{1}^{*} \\
\vdots & \vdots & \vdots & \vdots \\
0 & 0 & c_{-n+1}^{*}+h_{n-3}^{*} & c_{-n+2}^{*}+h_{n-2}^{*} \\
0 & 0 & h_{n-2}^{*} & c_{-n+1}^{*}+h_{n-1}^{*} \\
0 & 0 & h_{n-1}^{*} & h_{n}^{*} \\
\vdots & \vdots & \vdots & \vdots \\
0 & 0 & h_{2(n-2)}^{*} & h_{2(n-1)}^{*}
\end{array}\right] .
$$

Example 2 (Displacement rank 5). Define the quantities

$$
\begin{aligned}
& {\left[\begin{array}{lllll}
\gamma_{i}^{*} & -\delta_{i}^{*} & \eta_{i}^{*} & -\kappa_{i}^{*} & \xi_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} G_{i}(0),} \\
& {\left[\begin{array}{lllll}
\mu_{i}^{*} & -\sigma_{i}^{*} & \chi_{i}^{*} & -\psi_{i}^{*} & \varphi_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} B_{i}(0),} \\
& U_{i}=\left[\begin{array}{ccccc}
\gamma_{i}^{-*} & \gamma_{i}^{-*} \delta_{i}^{*} & -\gamma_{i}^{-*} \eta_{i}^{*} & \gamma_{i}^{-*} \kappa_{i}^{*} & \varphi_{i} \\
0 & 1 & 0 & 0 & \psi_{i} \\
0 & 0 & 1 & 0 & \chi_{i} \\
0 & 0 & 0 & 1 & \sigma_{i} \\
0 & 0 & 0 & 0 & \mu_{i}
\end{array}\right], \\
& V_{i}=\left[\begin{array}{ccccc}
\mu_{i}^{*} & \mu_{i}^{-*} \sigma_{i}^{*} & -\mu_{i}^{-*} \chi_{i}^{*} & \mu_{i}^{-*} \psi_{i}^{*} & \xi_{i} \\
0 & 1 & 0 & 0 & \kappa_{i} \\
0 & 0 & 1 & 0 & \eta_{i} \\
0 & 0 & 0 & 1 & \delta_{i} \\
0 & 0 & 0 & 0 & \gamma_{i}
\end{array}\right] .
\end{aligned}
$$

From (18), it follows that

$$
\begin{equation*}
-\xi_{i}^{*} \mu_{i}+\kappa_{i}^{*} \sigma_{i}-\eta_{i}^{*} \chi_{i}+\delta_{i}^{*} \psi_{i}-\gamma_{i}^{*} \varphi_{i}=0 \tag{26}
\end{equation*}
$$

It is easy to check with the help of (26) that (19)-(21) are satisfied. Observe that $U_{i}$ and $V_{i}$ are sparse matrices. This leads to an overall $O\left(4 \rho n^{2}\right)$ computational complexity.

## 6. FACTORIZATION OF QT AND QH MATRICES

Algorithms for the triangular factorization of quasi-Toeplitz matrices (QT) and quasi-Hankel matrices ( QH ) follow readily as special cases of the general recursive procedure described by Theorem 1.

### 6.1. Quasi-Hankel Matrices

According to Equations (1) and (3) the generating function of a quasiHankel matrix can be expressed in the form

$$
\mathrm{QH}(z, w)=\frac{G(z) J B^{*}(w)}{j\left(z-w^{*}\right)}
$$

where

$$
J=\left[\begin{array}{cc}
0 & -j \\
j & 0
\end{array}\right]
$$

$$
G(z)=\left[\begin{array}{ll}
\alpha(z) & z \alpha(z) h(z)
\end{array}\right], \text { and } B(z)=\left[\begin{array}{ll}
\beta(z) & z \beta(z) \tilde{h}(z)
\end{array}\right]
$$

Notice that $d(z, w)=j\left(z-w^{*}\right)$ can be written as

$$
d(z, w)=\left[\begin{array}{ll}
1 & z
\end{array}\right] J\left[\begin{array}{c}
1 \\
w^{*}
\end{array}\right],
$$

which shows that $d(z, w)$ satisfies (15), since $J$ has one positive and one negative eigenvalue. Following steps similar to those at the beginning of Section 5, we get

$$
\lambda(z)=j\left(\frac{1}{\tau^{*}}-\frac{1}{z}\right)
$$

which simplifies to $\lambda(z)=-j z^{-1}$ on setting $\tau=\infty$ [note that any choice of $\tau$ with $\operatorname{Im} \tau=0$ satisfies $d(\tau, \tau)=0$ ]. Now define

Then it is easy to check that

$$
\begin{aligned}
U_{i} J V_{i}^{*} & =J \\
G_{i}(0) U_{i} & =\left|Q H_{i}(0,0)\right|^{1 / 2}\left[\begin{array}{ll}
1 & 0
\end{array}\right], \\
B_{i}(0) V_{i} & =\left|Q H_{i}(0,0)\right|^{1 / 2}\left[\begin{array}{ll}
1 & 0
\end{array}\right] .
\end{aligned}
$$

$$
\begin{aligned}
& Q H_{i}(z, w)=\frac{G_{i}(z) J B_{i}^{*}(w)}{j\left(z-w^{*}\right)}, \\
& {\left[\begin{array}{ll}
\gamma_{i}^{*} & -\delta_{i}^{*}
\end{array}\right]=\left|Q H_{i}(0,0)\right|^{-1 / 2} G_{i}(0),} \\
& {\left[\begin{array}{ll}
\mu_{i}^{*} & \left.-\sigma_{i}^{*}\right]=\left|Q H_{i}(0,0)\right|^{-1 / 2} B_{i}(0), ~
\end{array}\right.} \\
& U_{i}=\left[\begin{array}{cc}
\gamma_{i}^{-*} & \sigma_{i} \\
0 & \mu_{i}
\end{array}\right], \quad \text { and } \quad V_{i}=\left[\begin{array}{cc}
\mu_{i}^{-*} & \delta_{i} \\
0 & \gamma_{i}
\end{array}\right] \text {. }
\end{aligned}
$$

and thus we get the following result (which can be written in matrix form as well).

ALGORITHM 2 (Factorization of QH matrices). The triangular factorization of a strongly regular quasi-Hankel matrix QH can be carried out by the following recursive procedure:

$$
\begin{aligned}
z G_{i+1}(z) & =G_{i}(z)\left[\begin{array}{cc}
\gamma_{i}^{-*} & \sigma_{i} \\
0 & \mu_{i}
\end{array}\right]\left[\begin{array}{cc}
1 & 0 \\
-\epsilon_{i} z^{-1} & 1
\end{array}\right], \\
z B_{i+1}(z) & =B_{i}(z)\left[\begin{array}{cc}
\mu_{i}^{-*} & \delta_{i} \\
0 & \gamma_{i}
\end{array}\right]\left[\begin{array}{cc}
1 & 0 \\
-\epsilon_{i}^{*} z^{-1} & 1
\end{array}\right], \\
d_{i} & =G_{i}^{\prime}(0) \tilde{J} B_{i}^{*}(0), \quad \epsilon_{i}=\frac{\left|d_{i}\right|}{d_{i}} \\
l_{i}(z) & =z^{i-1} G_{i}(z) \tilde{J} B_{i}^{*}(0) d_{i}^{-1}, \quad \text { and } \\
u_{i}(w) & =-\left(w^{*}\right)^{i-1} G_{i}(0) \tilde{J B_{i}^{*}}(w) d_{i}^{-1}
\end{aligned}
$$

where $J=j \tilde{J}$. Note that when QH is Hermitian we get $G_{i}(z)=B_{i}(z)$, $U_{i}=V_{i}$, and the previous recursions reduce to those given in [23].

### 6.2. Quasi-Toeplitz Matrices

The generating function of a quasi-Toeplitz matrix can be expressed in the form

$$
\operatorname{QT}(z, w)=\frac{G(z) J B^{*}(w)}{1-z w^{*}}
$$

where $G(z)$ and $B(z)$ are $1 \times 2$ row vector functions, and

$$
J=\left[\begin{array}{rr}
1 & 0 \\
0 & -1
\end{array}\right]
$$

Notice that $d(z, w)=1-z w^{*}$ can be written as

$$
d(z, w)=\left[\begin{array}{ll}
1 & z
\end{array}\right] J\left[\begin{array}{c}
1 \\
w^{*}
\end{array}\right]
$$

which shows that $d(z, w)$ satisfies (15), since $J$ has one positive and one negative eigenvalue. Following steps similar to those at the beginning of Section 5, we get

$$
\lambda(z)=1-z \tau^{*}
$$

which simplifies to $\lambda(z)=1-z$ on setting $\tau=1$ [note that any choice of $\tau$ with $|\tau|=1$ satisfies $d(\tau, \tau)=0$ ]. Now define

$$
\begin{aligned}
& \mathrm{QT}_{i}(z, w)=\frac{G_{i}(z) J B_{i}^{*}(w)}{1-z w^{*}}, \\
& G_{i}(z)=\left[\begin{array}{ll}
s_{i}(z) & v_{i}(z)
\end{array}\right]=\left[\begin{array}{llll}
1 & z & z^{2} & \ldots
\end{array}\right]\left[\begin{array}{cc}
s_{i i} & v_{i i} \\
s_{i+1, i} & v_{i+1, i} \\
\vdots & \vdots
\end{array}\right], \\
& B_{i}(z)=\left[\begin{array}{ll}
x_{i}(z) & y_{i}(z)
\end{array}\right]=\left[\begin{array}{llll}
1 & z & z^{2} & \ldots
\end{array}\right]\left[\begin{array}{cc}
x_{i i} & y_{i i} \\
x_{i+1, i} & y_{i+1, i} \\
\vdots & \vdots
\end{array}\right] \text {, } \\
& k_{i}=v_{i i} / s_{i i} \quad \text { and } \quad \xi_{i}=y_{i i} / x_{i i} \text {, } \\
& U_{i}=\frac{1}{1-k_{i} \xi_{i}^{*}}\left[\begin{array}{cc}
1 & -k_{i} \\
-\xi_{i}^{*} & 1
\end{array}\right] \quad \text { and } \quad V_{i}=\left[\begin{array}{cc}
1 & -\xi_{i} \\
-k_{i}^{*} & 1
\end{array}\right] .
\end{aligned}
$$

Then it is easy to check that

$$
\begin{aligned}
U_{i} J V_{i}^{*} & =J \\
G_{i}(0) U_{i} & =s_{i i}\left[\begin{array}{ll}
1 & 0
\end{array}\right], \\
B_{i}(0) V_{i} & =x_{i i}\left(1-k_{i}^{*} \xi_{i}\right)\left[\begin{array}{ll}
1 & 0
\end{array}\right], \\
s_{i i} x_{i i}^{*}\left(1-k_{i} \xi_{i}^{*}\right) & =Q T_{i}(0,0),
\end{aligned}
$$

which yields the following result.

ALGORITHM 3 (Factorization of QT matrices). The triangular factorization of a strongly regular quasi-Toeplitz matrix QT can be carried out by the following recursive procedure:

$$
\begin{gathered}
z G_{i+1}(z)=\frac{1}{1-k_{i} \xi_{i}^{*}} G_{i}(z)\left[\begin{array}{cc}
1 & -k_{i} \\
-\xi_{i}^{*} & 1
\end{array}\right]\left[\begin{array}{cc}
z & 0 \\
0 & 1
\end{array}\right] \\
z B_{i+1}(z)=B_{i}(z)\left[\begin{array}{cc}
1 & -\xi_{i} \\
-k_{i}^{*} & 1
\end{array}\right]\left[\begin{array}{cc}
z & 0 \\
0 & 1
\end{array}\right] \\
d_{i}=G_{i}(0) J B_{i}^{*}(0) \\
l_{i}(z)=z^{i} G_{i}(z) J B_{i}^{*}(0) d_{i}^{-1} \text { and } u_{i}(w)=\left(w^{*}\right)^{i} G_{i}(0) J B_{i}^{*}(w) d_{i}^{-1} \\
k_{i}=v_{i i} / s_{i i} \text { and } \xi_{i}=y_{i i} / x_{i i} .
\end{gathered}
$$

Note that when QT is Hermitian positive definite we get $G_{i}(z)=B_{i}(z)$, $k_{i}=\xi_{i}$, and the previous recursions reduce to

$$
\begin{align*}
z \mathrm{G}_{i+1}(z) & =\mathrm{G}_{i}(z) \frac{1}{\sqrt{1-\left|k_{i}\right|^{2}}}\left[\begin{array}{cc}
1 & -k_{i} \\
-k_{i}^{*} & 1
\end{array}\right]\left[\begin{array}{cc}
z & 0 \\
0 & 1
\end{array}\right]  \tag{27}\\
k_{i} & =\boldsymbol{v}_{i i} / s_{i i}
\end{align*}
$$

which is the linearized form of the conventional Schur recursion [23, 27, 28]. If we define $\bar{v}_{i}(z)=z^{i} v_{i}(z), \bar{s}_{i}(z)=z^{i} s_{i}(z)$, and $f_{i}(z)=\bar{v}_{i}(z) / \bar{s}_{i}(z)$, then (27) reduces to the classical Schur recursion [29] for functions that are analytic and bounded by unity in the unit dise,

$$
f_{i+1}(z)=\frac{1}{z} \frac{f_{i}(z)-k_{i}}{1-k_{i}^{*} f_{i}(z)} \quad \text { with } \quad k_{i}=f_{i}(0)
$$

The positive definiteness of QT guarantees $\left|k_{i}\right|<1$.
The recursions for the non-Hermitian quasi-Toeplitz case are the same expressions derived in [30], though from a different point of view. We remark that the recursions for $G_{i}(z)$ and $B_{i}(z)$ differ by a scaling factor. One can, however, choose different matrices $U_{i}$ and $V_{i}$ so that the resulting recursions for the generators $G_{i}$ and $B_{i}$ become similar.

## 7. CONCLUSION

In this paper we have presented fast triangular factorization algorithms [ $O\left(n^{2}\right)$ ] for matrices close to Toeplitz plus Hankel, and in particular for quasi-Hankel, quasi-Toeplitz, and Toeplitz-plus-Hankel matrices. Both Hermitian and non-Hermitian matrices were considered. The arguments are based on an extension to the non-Hermitian case of the generating-function approach discussed in [23]. The results include many special cases studied separately in the literature. We remark that we can, as well, compute the triangular factorization of the inverse of an $n \times n$ close to Toeplitz-plusHankel matrix $R$. For this purpose, we define the extended matrix

$$
\left[\begin{array}{cc}
-R & I \\
I & 0
\end{array}\right]
$$

and observe that its Schur complement with respect to the ( 1,1 ) block entry is $R^{-1}$. We can exploit the structure of this extended matrix in the derivation of a fast factorization algorithm for $R^{-1}$; similar ideas can be used for solving linear equations with $R$ as the coefficient matrix, finding orthogonal factorizations, etc. See [31] for details.

## APPENDIX A

We list here possible choices of $U_{i}$ and $V_{i}$ for other displacement ranks than the cases $\rho=4$ and $\rho=5$ treated in Section 5.

Displacement rank 2:

$$
\begin{aligned}
& {\left[\begin{array}{ll}
\gamma_{i}^{*} & -\delta_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} G_{i}(0),} \\
& {\left[\begin{array}{ll}
\mu_{i}^{*} & -\sigma_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} B_{i}(0),} \\
& U_{i}=\left[\begin{array}{cc}
\gamma_{i}^{-*} & \sigma_{i} \\
0 & \mu_{i}
\end{array}\right], \quad \text { and } \quad V_{i}=\left[\begin{array}{cc}
\mu_{i}^{-*} & \delta_{i} \\
0 & \gamma_{i}
\end{array}\right] .
\end{aligned}
$$

Displacement rank 3:

$$
\begin{gathered}
{\left[\begin{array}{lll}
\gamma_{i}^{*} & -\delta_{i}^{*} & \eta_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} G_{i}(0),} \\
{\left[\begin{array}{lll}
\mu_{i}^{*} & -\sigma_{i}^{*} & \chi_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} B_{i}(0),} \\
U_{i}=\left[\begin{array}{ccc}
\gamma_{i}^{-*} & \gamma_{i}^{-*} \delta_{i}^{*} & \chi_{i} \\
0 & 1 & \sigma_{i} \\
0 & 0 & \mu_{i}
\end{array}\right], \quad \text { and } V_{i}=\left[\begin{array}{ccc}
\mu_{i}^{-*} & \mu_{i}^{-*} \sigma_{i}^{*} & \eta_{i} \\
0 & 1 & \delta_{i} \\
0 & 0 & \gamma_{i}
\end{array}\right] .
\end{gathered}
$$

Displacement rank 6:

$$
\left.\begin{array}{rl}
{\left[\begin{array}{cccccc}
\gamma_{i}^{*} & -\delta_{i}^{*} & \eta_{i}^{*} & -\kappa_{i}^{*} & \xi_{i}^{*} & -\vartheta_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} G_{i}(0),} \\
\mu_{i}^{*} & -\sigma_{i}^{*} \\
\chi_{i}^{*} & -\psi_{i}^{*} \\
U_{i} & \varphi_{i}^{*} \\
U_{i} & -\theta_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} B_{i}(0), ~\left[\begin{array}{ccccccc}
\gamma_{i}^{-*} & \delta_{i}^{*} & \delta_{i}^{*}-\frac{\eta_{i}^{*}}{\mu_{i} \gamma_{i}^{*}} & \kappa_{i}^{*}-\frac{\eta_{i}^{*}}{\mu_{i} \gamma_{i}^{*}} & \delta_{i}^{*}+\kappa_{i}^{*}-\frac{\xi_{i}^{*}+\eta_{i}^{*}}{\mu_{i} \gamma_{i}^{*}} & \theta_{i} \\
0 & \gamma_{i}^{*} & \gamma_{i}^{*} & 0 & \gamma_{i}^{*} & \varphi_{i} \\
0 & 0 & \mu_{i}^{-1} & \mu_{i}^{-1} & \mu_{i}^{-1} & \psi_{i} \\
0 & 0 & 0 & \gamma_{i}^{*} & \gamma_{i}^{*} & \chi_{i} \\
0 & 0 & 0 & 0 & \mu_{i}^{-1} & \sigma_{i} \\
0 & 0 & 0 & 0 & 0 & \mu_{i}
\end{array}\right],
$$

Displacement rank 7:

$$
\begin{aligned}
& {\left[\begin{array}{lllllll}
\gamma_{i}^{*} & -\delta_{i}^{*} & \eta_{i}^{*} & -\kappa_{i}^{*} & \xi_{i}^{*} & -\vartheta_{i}^{*} & \nu_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} G_{i}(0)} \\
& {\left[\begin{array}{lllllll}
\mu_{i}^{*} & -\sigma_{i}^{*} & \chi_{i}^{*} & -\psi_{i}^{*} & \varphi_{i}^{*} & -\theta_{i}^{*} & \mathbf{s}_{i}^{*}
\end{array}\right]=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} B_{i}(0)}
\end{aligned}
$$

$$
U_{i}=\left[\begin{array}{ccccccc}
\gamma_{i}^{-*} & \gamma_{i}^{-*} \delta_{i}^{*} & -\gamma_{i}^{-*} \boldsymbol{\eta}_{i}^{*} & \gamma_{i}^{-*} \boldsymbol{\kappa}_{i}^{*} & -\gamma_{i}^{-*} \xi_{i}^{*} & \gamma_{i}^{-*} \boldsymbol{\vartheta}_{i}^{*} & \boldsymbol{s}_{i} \\
0 & 1 & 0 & 0 & 0 & 0 & \theta_{i} \\
0 & 0 & 1 & 0 & 0 & 0 & \varphi_{i} \\
0 & 0 & 0 & 1 & 0 & 0 & \psi_{i} \\
0 & 0 & 0 & 0 & 1 & 0 & \chi_{i} \\
0 & 0 & 0 & 0 & 0 & 1 & \sigma_{i} \\
0 & 0 & 0 & 0 & 0 & 0 & \mu_{i}
\end{array}\right]
$$

$$
V_{i}=\left[\begin{array}{ccccccc}
\mu_{i}^{-*} & \mu_{i}^{-*} \sigma_{i}^{*} & -\mu_{i}^{-*} \chi_{i}^{*} & \mu_{i}^{-*} \psi_{i}^{*} & -\mu_{i}^{-*} \varphi_{i}^{*} & \mu_{i}^{-*} \theta_{i}^{*} & \nu_{i} \\
0 & 1 & 0 & 0 & 0 & 0 & \boldsymbol{\vartheta}_{i} \\
0 & 0 & 1 & 0 & 0 & 0 & \xi_{i} \\
0 & 0 & 0 & 1 & 0 & 0 & \kappa_{i} \\
0 & 0 & 0 & 0 & 1 & 0 & \eta_{i} \\
0 & 0 & 0 & 0 & 0 & 1 & \delta_{i} \\
0 & 0 & 0 & 0 & 0 & 0 & \gamma_{i}
\end{array}\right] .
$$

Displacement rank 8:

$$
\begin{aligned}
& {\left[\begin{array}{llllllll}
\gamma_{i}^{*} & -\delta_{i}^{*} & \eta_{i}^{*} & -\kappa_{i}^{*} & \xi_{i}^{*} & -\vartheta_{i}^{*} & \nu_{i}^{*} & -\varrho_{i}^{*}
\end{array}\right]} \\
& \\
& \quad=\left|\tilde{R}_{i}(0,0)\right|^{-1 / 2} G_{i}(0) \\
& {\left[\begin{array}{lllllll}
\mu_{i}^{*} & -\sigma_{i}^{*} & \chi_{i}^{*} & -\psi_{i}^{*} & \varphi_{i}^{*} & -\theta_{i}^{*} & \varsigma_{i}^{*} \\
& -\varpi_{i}^{*}
\end{array}\right]} \\
& \\
& =
\end{aligned}
$$

$$
\begin{aligned}
& U_{i}=\left[\begin{array}{cccccccc}
\gamma_{i}^{-*} & \delta_{i}^{*} & \delta_{i}^{*}-\frac{\eta_{i}^{*}}{\mu_{i} \gamma_{i}^{*}} & \kappa_{i}^{*}-\eta_{i}^{*} & u_{i 1} & u_{i 2} & u_{i 3} & \varpi_{i} \\
0 & \gamma_{i}^{*} & \gamma_{i}^{*} & 0 & \gamma_{i}^{*} & \gamma_{i}^{*} & \gamma_{i}^{*} & s_{i} \\
0 & 0 & \mu_{i}^{-1} & \gamma_{i}^{*} & \gamma_{i}^{*} & \mu_{i}^{-1} & \gamma_{i}^{*} & \theta_{i} \\
0 & 0 & 0 & \gamma_{i}^{*} & \gamma_{i}^{*} & 0 & \gamma_{i}^{*} & \varphi_{i} \\
0 & 0 & 0 & 0 & \mu_{i}^{-1} & \gamma_{i}^{*} & \gamma_{i}^{*} & \psi_{i} \\
0 & 0 & 0 & 0 & 0 & \gamma_{i}^{*} & \gamma_{i}^{*} & \chi_{i} \\
0 & 0 & 0 & 0 & 0 & 0 & \mu_{i}^{-1} & \sigma_{i} \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & \mu_{i}
\end{array}\right], \\
& V_{i}=\left[\begin{array}{cccccccc}
\mu_{i}^{-*} & \sigma_{i}^{*} & \sigma_{i}^{*}-\frac{\chi_{i}^{*}}{\gamma_{i} \mu_{i}^{*}} & \psi_{i}^{*}-\xi_{i}^{*} & v_{i 1} & \boldsymbol{v}_{i 2} & v_{i 3} & \varrho_{i} \\
0 & \mu_{i}^{*} & \mu_{i}^{*} & 0 & \mu_{i}^{*} & \mu_{i}^{*} & \mu_{i}^{*} & \nu_{i} \\
0 & 0 & \gamma_{i}^{-1} & \mu_{i}^{*} & \mu_{i}^{*} & \gamma_{i}^{-1} & \mu_{i}^{*} & \vartheta_{i} \\
0 & 0 & 0 & \mu_{i}^{*} & \mu_{i}^{*} & 0 & \mu_{i}^{*} & \xi_{i} \\
0 & 0 & 0 & 0 & \gamma_{i}^{-1} & \mu_{i}^{*} & \mu_{i}^{*} & \kappa_{i} \\
0 & 0 & 0 & 0 & 0 & \mu_{i}^{*} & \mu_{i}^{*} & \eta_{i} \\
0 & 0 & 0 & 0 & 0 & 0 & \gamma_{i}^{-1} & \delta_{i} \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & \gamma_{i}
\end{array}\right],
\end{aligned}
$$

where

$$
\begin{gathered}
u_{i 1}=\delta_{i}^{*}+\kappa_{i}^{*}-\eta_{i}^{*}-\frac{\xi_{i}^{*}}{\mu_{i} \gamma_{i}^{*}}, \quad u_{i 2}=\vartheta_{i}^{*}+\delta_{i}^{*}-\xi_{i}^{*}-\frac{\eta_{i}^{*}}{\mu_{i} \gamma_{i}^{*}} \\
u_{i 3}=\vartheta_{i}^{*}+\delta_{i}^{*}-\xi_{i}^{*}+\kappa_{i}^{*}-\eta_{i}^{*}-\frac{\nu_{i}^{*}}{\mu_{i} \gamma_{i}^{*}}, \\
v_{i 1}=\sigma_{i}^{*}+\psi_{i}^{*}-\chi_{i}^{*}-\frac{\varphi_{i}^{*}}{\gamma_{i} \mu_{i}^{*}}, \quad v_{i 2}=\theta_{i}^{*}+\sigma_{i}^{*}-\varphi_{i}^{*}-\frac{\chi_{i}^{*}}{\gamma_{i} \mu_{i}^{*}}, \\
v_{i 3}=\theta_{i}^{*}+\sigma_{i}^{*}-\varphi_{i}^{*}+\psi_{i}^{*}-\chi_{i}^{*}-\frac{\varsigma_{i}^{*}}{\gamma_{i} \mu_{i}^{*}} .
\end{gathered}
$$

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